

Thermal Charm production at LHC

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- Introduction
- Charm production
 - in QCD
 - in QGP
 - at LHC
 - from three-gluon interaction
- Summary

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Why is understanding charm production important in HIC?

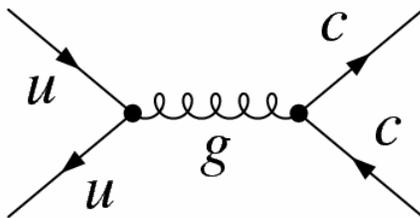
- Charmonium production: Braun-Munzinger, Thews, Greco
 - Yield depends quadratically on the charm quark number in statistical, kinetic, and coalescence models
 - Enhanced charm production would lead to possible charmonium enhancement instead of suppression, which was proposed as a signal for QGP (Matsui and Satz)
 - Expect charmonium suppression at RHIC but enhancement at LHC
- Charmed exotics production: Lee, Yasui, Liu & Ko (hep-ph/0707.1747)
 - Consideration of the color-spin interaction leads to possible stable charmed tetraquark meson T_{cc} ($ud \bar{c}\bar{c}$) and pentaquark baryon Θ_{sc} ($udus \bar{c}$)
 - Enhanced charm production at LHC makes the latter a possible factory for studying charmed exotics

Four stages of charm production in HIC

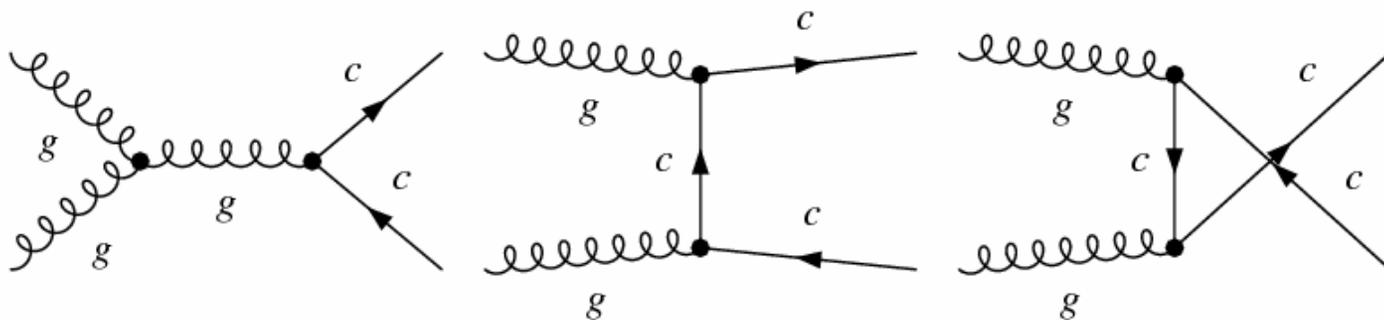
- Direct production: Mueller, Wang (92); Vogt (94); Gavin (96)
 - Mainly from initial gluon fusions
 - About 3 pairs in mid-rapidity at RHIC (from STAR collaboration)
 - About 20 pairs in mid-rapidity at LHC
- Pre-thermal production: Lin, Gyulassy (95), Levai, Mueller, Wang (95).....
 - Not important based on minijet gluons
 - Production from initial strong color field?
- Thermal production from QGP: Levai, Vogt (97)
 - Based on leading-order calculations
 - Important if initial temperature of QGP is high
- Thermal production from hadronic matter: Cassing et al. (99), Liu & Ko (02)
 - Such as $\pi N \rightarrow \Lambda_c D$ and $\rho N \rightarrow \Lambda_c D$
 - Expect small effect on charm production in HIC

Leading-order diagrams for charm production

1) $q\bar{q} \rightarrow c\bar{c}$

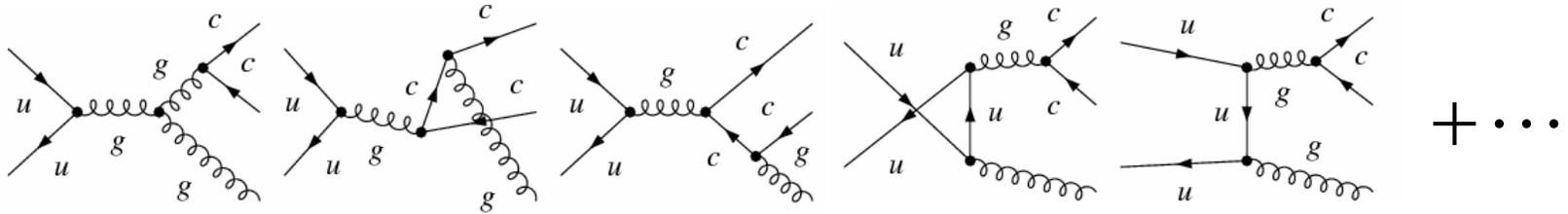


2) $gg \rightarrow c\bar{c}$

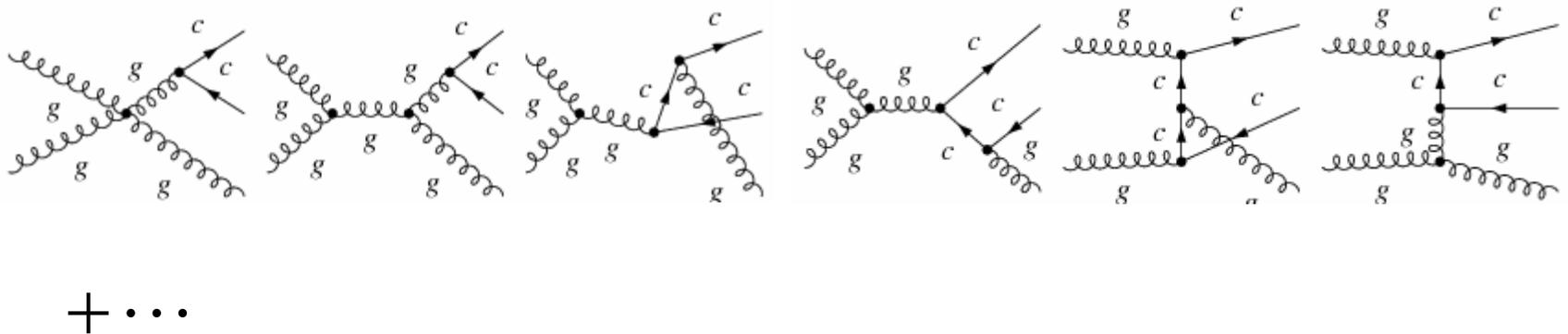


Next-Leading-order diagrams for charm production

1) $q\bar{q} \rightarrow c\bar{c}g$

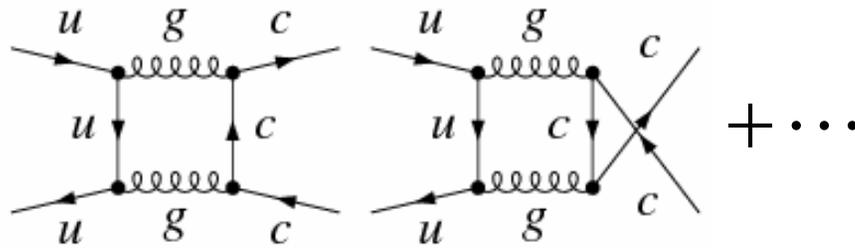


2) $gg \rightarrow c\bar{c}g$

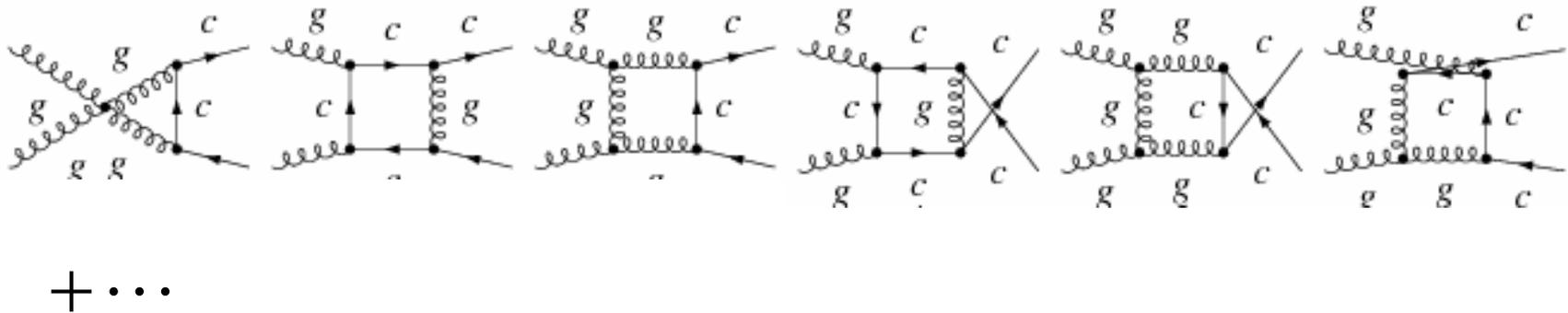


Virtual corrections to leading-order diagrams

1) $q\bar{q} \rightarrow c\bar{c}$

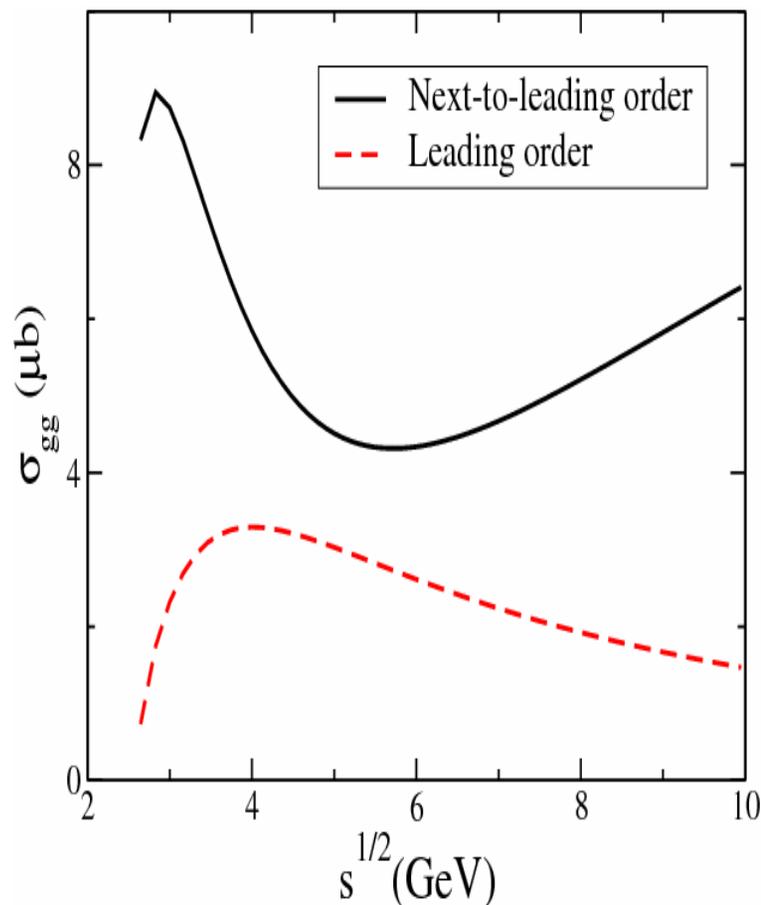
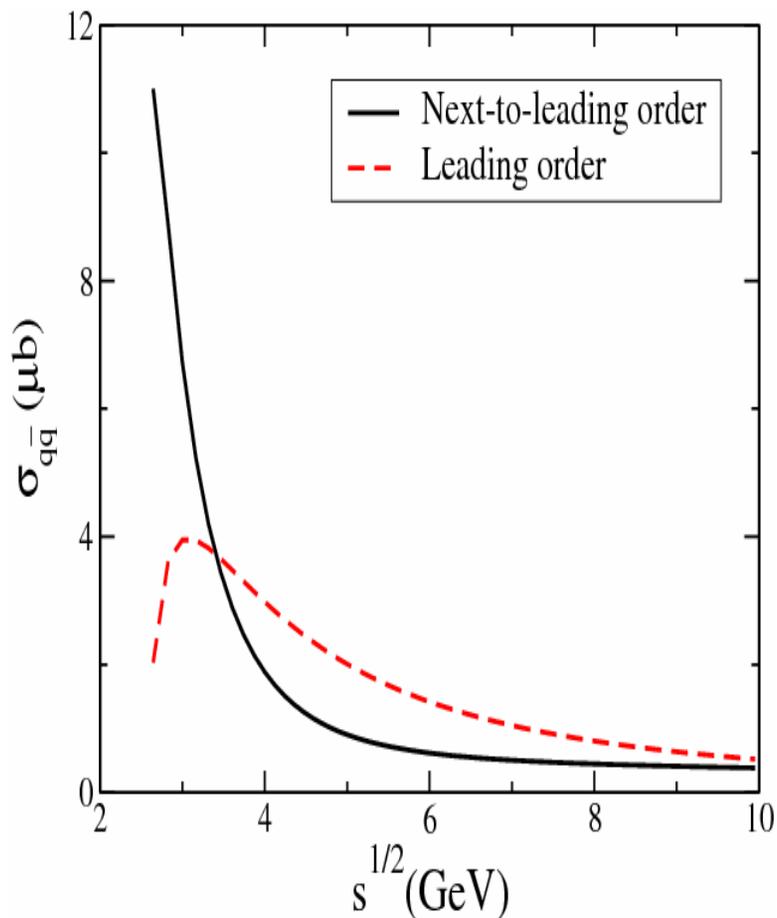


2) $gg \rightarrow c\bar{c}$



Charm quark production cross sections

P. Nason, S. Dawson & R.K. Ellis, NPB 303, 607 (1988)



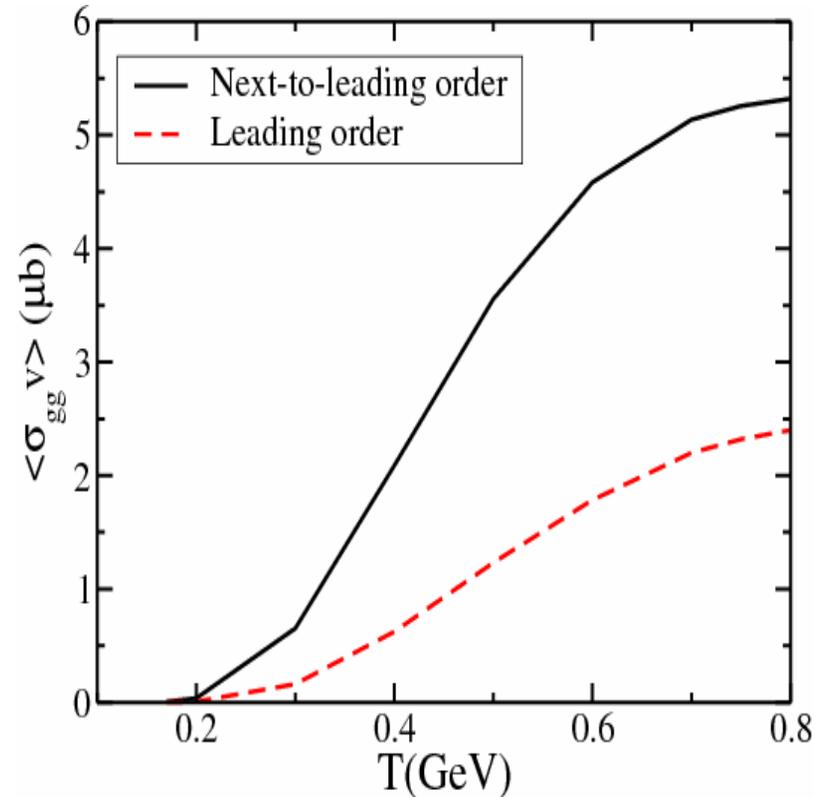
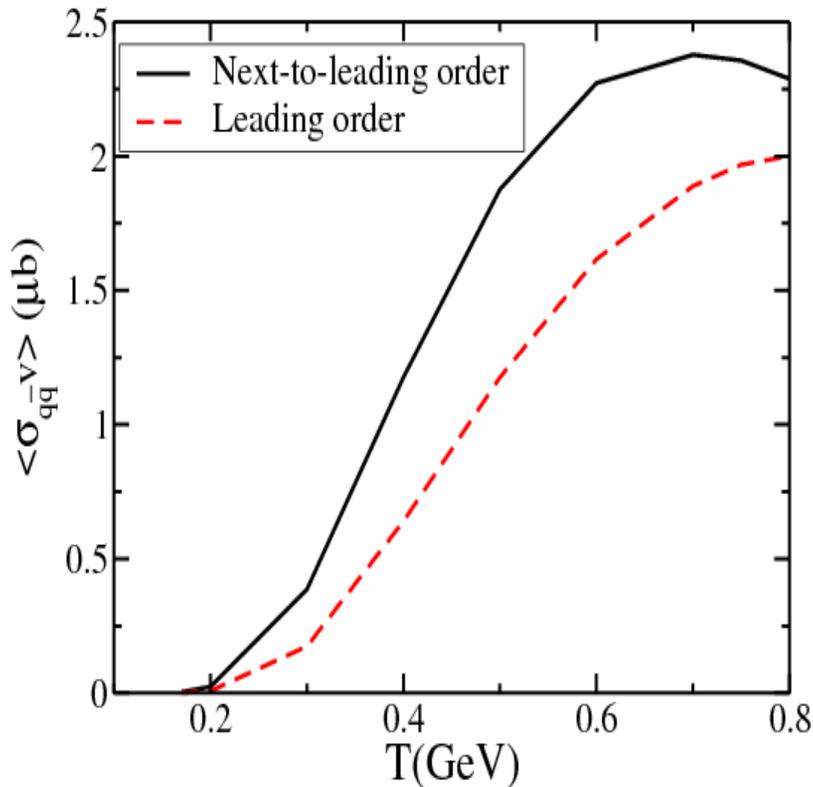
Next-to-leading order generally gives a larger cross section than the leading order except in $q\bar{q}$ annihilation at high energies.

Thermal averaged charm production cross sections

$$\langle \sigma_{ab \rightarrow cd} v \rangle = \frac{\int d^3 p_a d^3 p_b f_a(p_a) f_b(p_b) \sigma_{ab \rightarrow cd} v}{\int d^3 p_a d^3 p_b f_a(p_a) f_b(p_b)}$$

Thermal masses

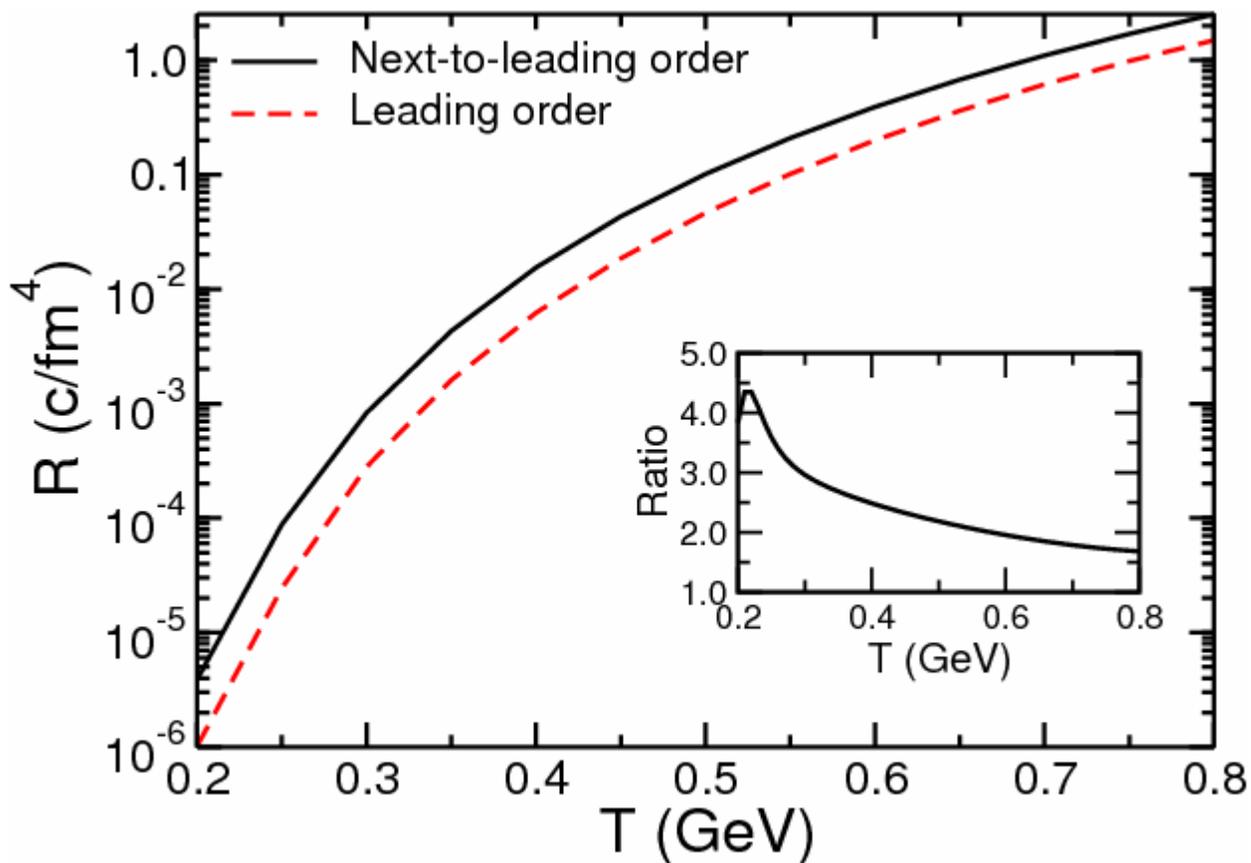
$$m_q = \frac{gT}{\sqrt{6}}, \quad m_g = \frac{gT}{\sqrt{2}}$$



Thermal averaged cross sections are larger in next-to-leading order, particularly in the gg channel. Slightly smaller if using massless partons.⁸

Charm production rate

$$R = \left[\langle \sigma_{q\bar{q} \rightarrow c\bar{c} \nu} \rangle + \langle \sigma_{q\bar{q} \rightarrow c\bar{c}g \nu} \rangle \right] n_q^{\text{eq}} n_{\bar{q}}^{\text{eq}} + \frac{1}{2} \left[\langle \sigma_{gg \rightarrow c\bar{c} \nu} \rangle + \langle \sigma_{gg \rightarrow c\bar{c}g \nu} \rangle \right] (n_g^{\text{eq}})^2$$



Production rate increases exponentially with temperature

Rate equation for charm production from QGP

$$\frac{1}{V} \frac{dN_{c\bar{c}}}{d\tau} \approx \left[\left(\langle \sigma_{q\bar{q} \rightarrow c\bar{c}} \rangle + \langle \sigma_{q\bar{q} \rightarrow c\bar{c}g} \rangle \right) n_q^{\text{eq}} n_{\bar{q}}^{\text{eq}} + \frac{1}{2} \left(\langle \sigma_{gg \rightarrow c\bar{c}} \rangle + \langle \sigma_{gg \rightarrow c\bar{c}g} \rangle \right) (n_g^{\text{eq}})^2 \right] \left[1 - \left(\frac{n_{c\bar{c}}}{n_{c\bar{c}}^{\text{eq}}} \right)^2 \right]$$

QGP fire-cylinder dynamics at LHC

- Longitudinally boost invariant and transversely accelerated → volume

$$V(\tau) = \pi \left[R_0 + \frac{a}{2} (\tau - \tau_0)^2 \right]^2 \tau$$

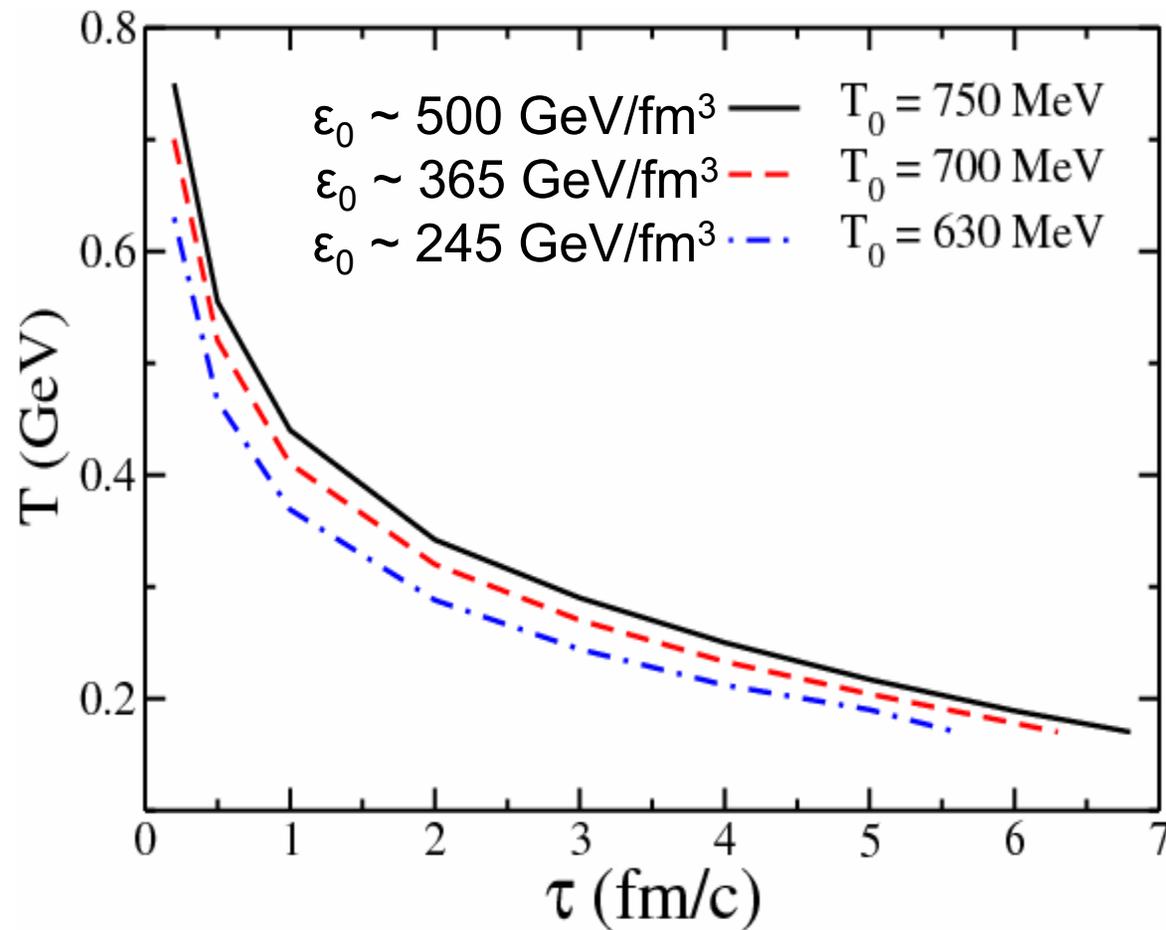
- For Pb+Pb @ 5.5 ATeV, $R_0 \sim 1.2 A^{1/3} \sim 7$ fm
- Expecting the QGP formation time τ_0 to be less than ~ 0.5 fm/c at RHIC, we take $\tau_0 = 0.2$ fm/c
- Taking transverse acceleration $a=0.1$ c²/fm, similar to that at RHIC

Initial temperature of QGP formed in HIC

- Color glass condensate: T. Lappi, PLB 643, 11 (2006)
 - At LHC, energy density at $\tau = 0.07$ fm/c: $\varepsilon \sim 700$ GeV/fm³
 - Assuming ε decreases with time as $1/\tau \rightarrow \varepsilon_0 \sim 245$ GeV/fm³ at $\tau_0 = 0.2$ fm/c
 - Using $\varepsilon \sim (T/160)^4$ GeV/fm³ $\rightarrow T_0 \sim 633$ MeV at LHC
 - At RHIC, $\varepsilon \sim 130$ GeV/fm³ at $\tau = 0.1$ fm/c $\rightarrow T_0 \sim 361$ MeV at $\tau_0 = 0.5$ fm/c
 - Uncertainty is , however, large due to Q_s^4 dependence
- HIJING (Gyulassy and Wang) or AMPT: Lin et al., PRC 72, 064901 (2005)
 - Initial transverse energy $dE_T/dy \sim 3000$ GeV at LHC
 - $$\varepsilon_0 \approx \frac{dE_T/dy}{\pi R_0^2 \tau_0} \approx \frac{3000}{3 \times 4.7^2 \times 0.2} \approx 226 \text{ GeV/fm}^3 \rightarrow T_0 \approx 620 \text{ MeV}$$
 - At RHIC, $dE_T/dy \sim 1000$ GeV $\rightarrow \varepsilon_0 \sim 33$ GeV/fm³ $\rightarrow T_0 \sim 383$ MeV at $\tau_0 = 0.5$ fm/c

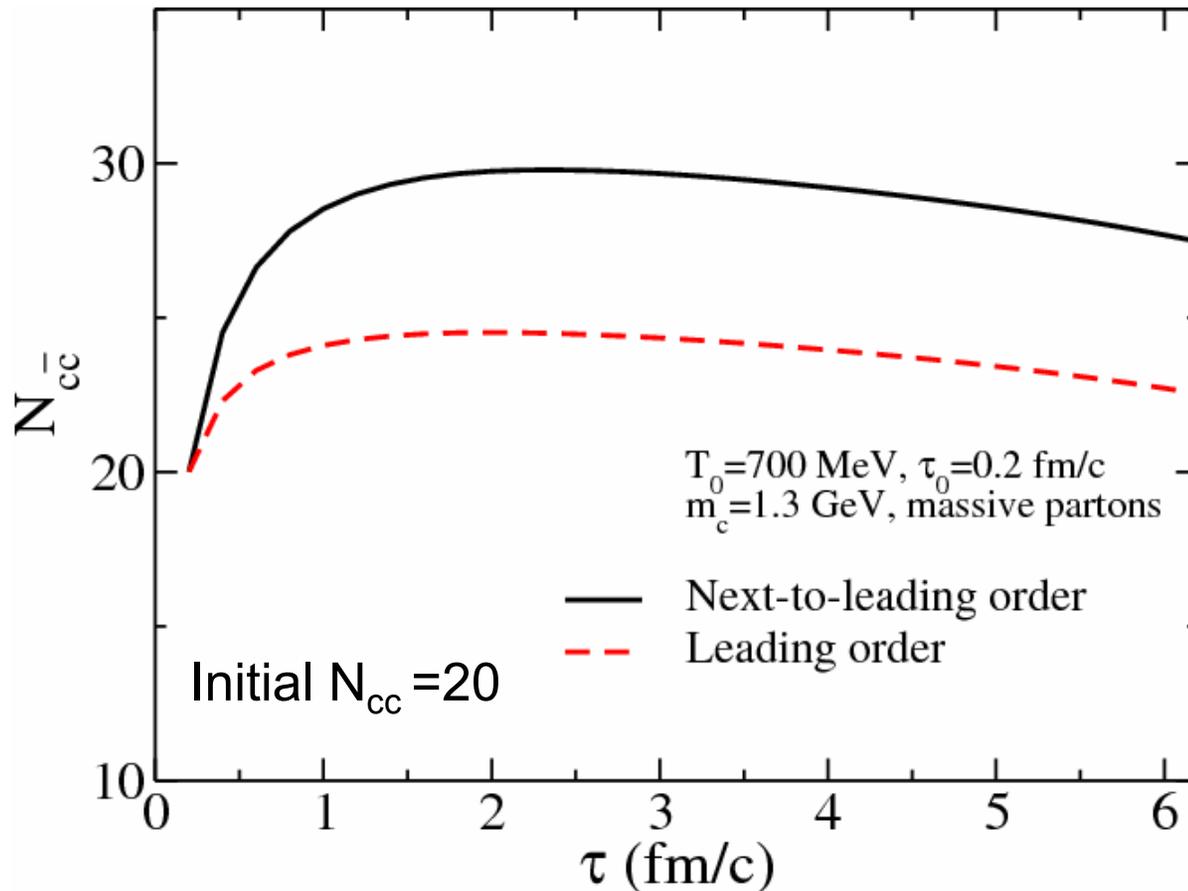
Temperature evolution at LHC

Entropy conservation \rightarrow



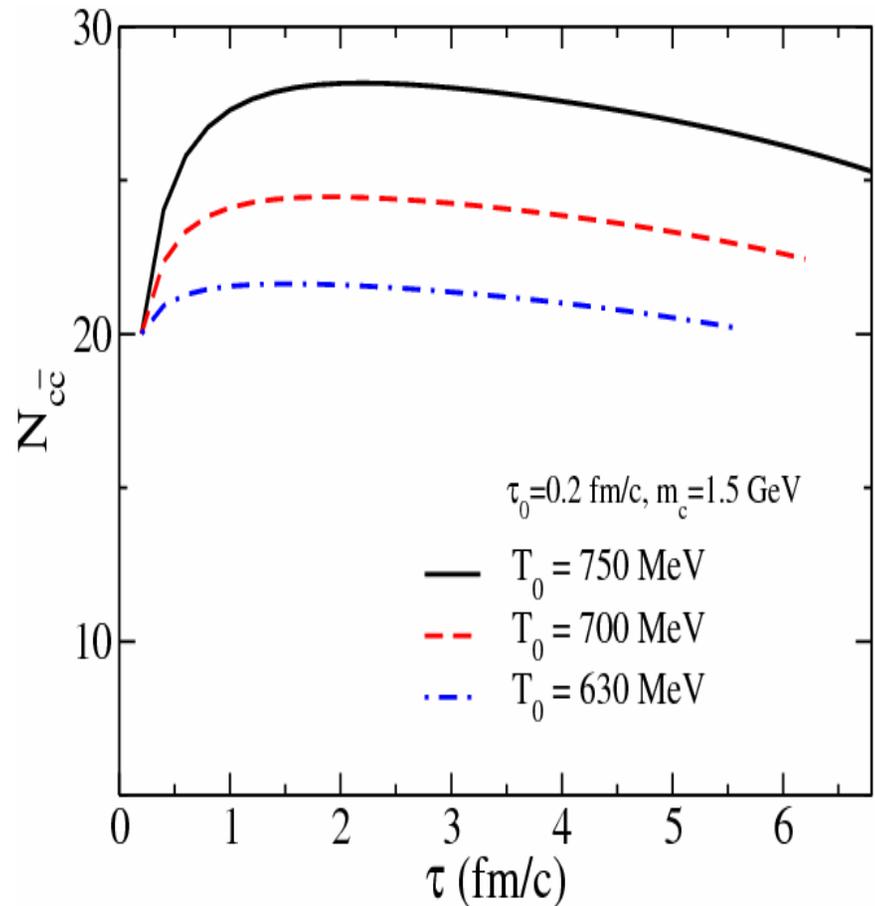
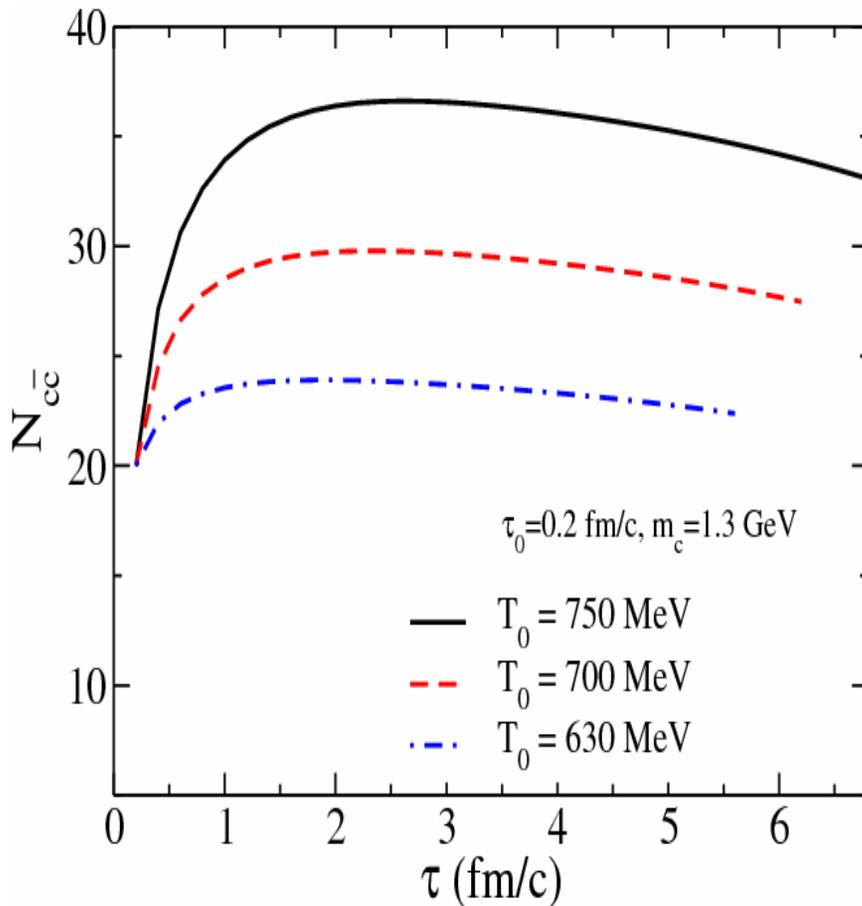
High temperature only exists briefly during early stage of QGP

Time evolution of charm quark pair at LHC



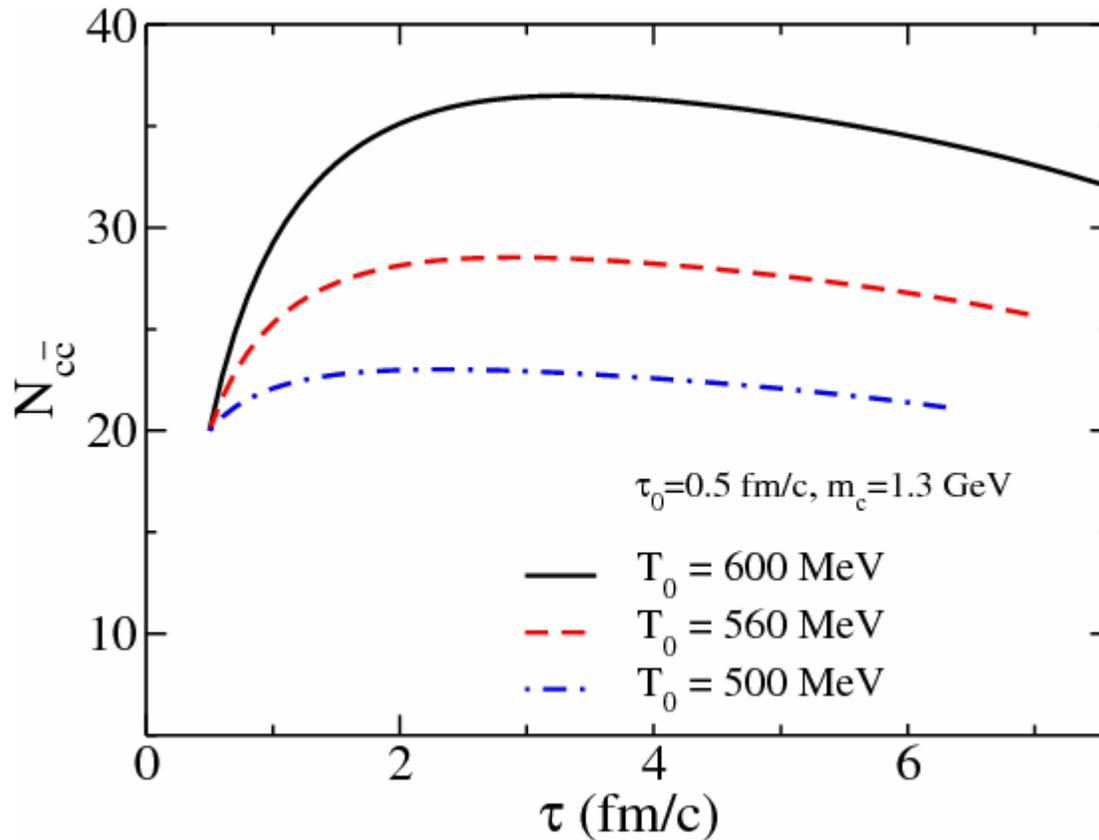
- Charm production in next-to-leading order is more than a factor of two larger than in the leading order
- Results using massless gluons are slightly larger

Initial temperature and charm quark mass dependence of thermal charm production



Increases with initial temperature but decreases with charm quark mass.

Charm production at LHC for $\tau_0=0.5$ fm/c

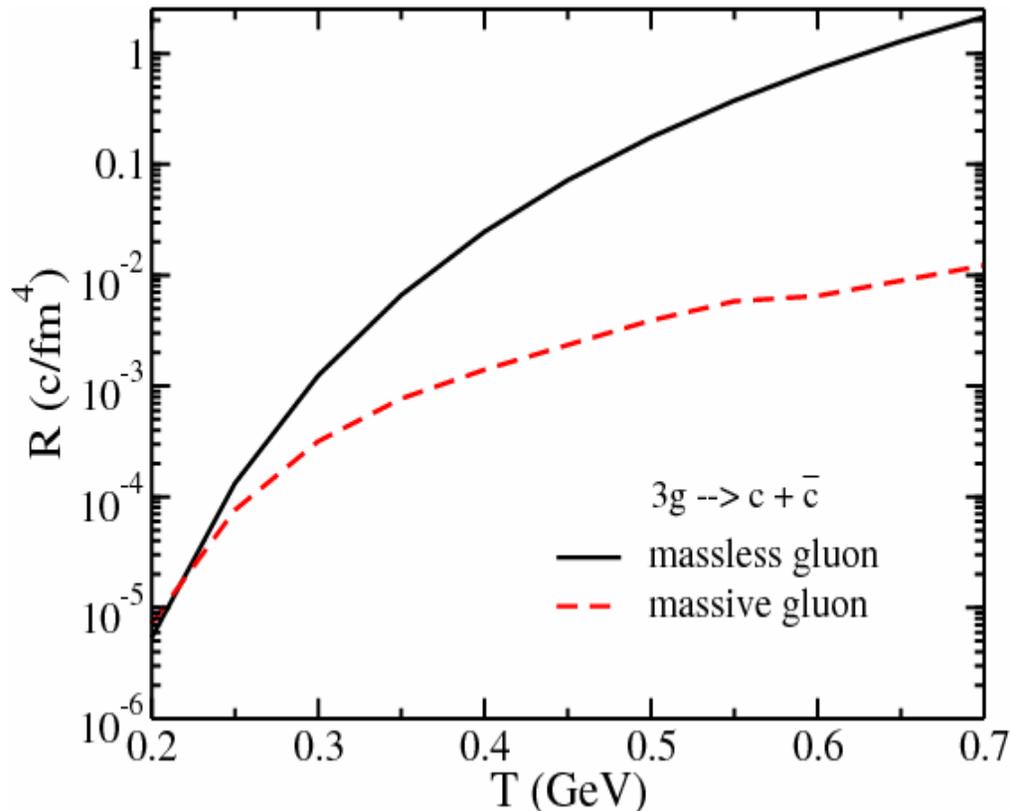


Similar results as $\tau_0 = 0.2$ fm/c, although initial temperature is lower

Charm production from three-gluon interaction $ggg \rightarrow c\bar{c}$

Determine rate for $ggg \rightarrow c\bar{c}$ from $c\bar{c} \rightarrow ggg$ via detailed balance

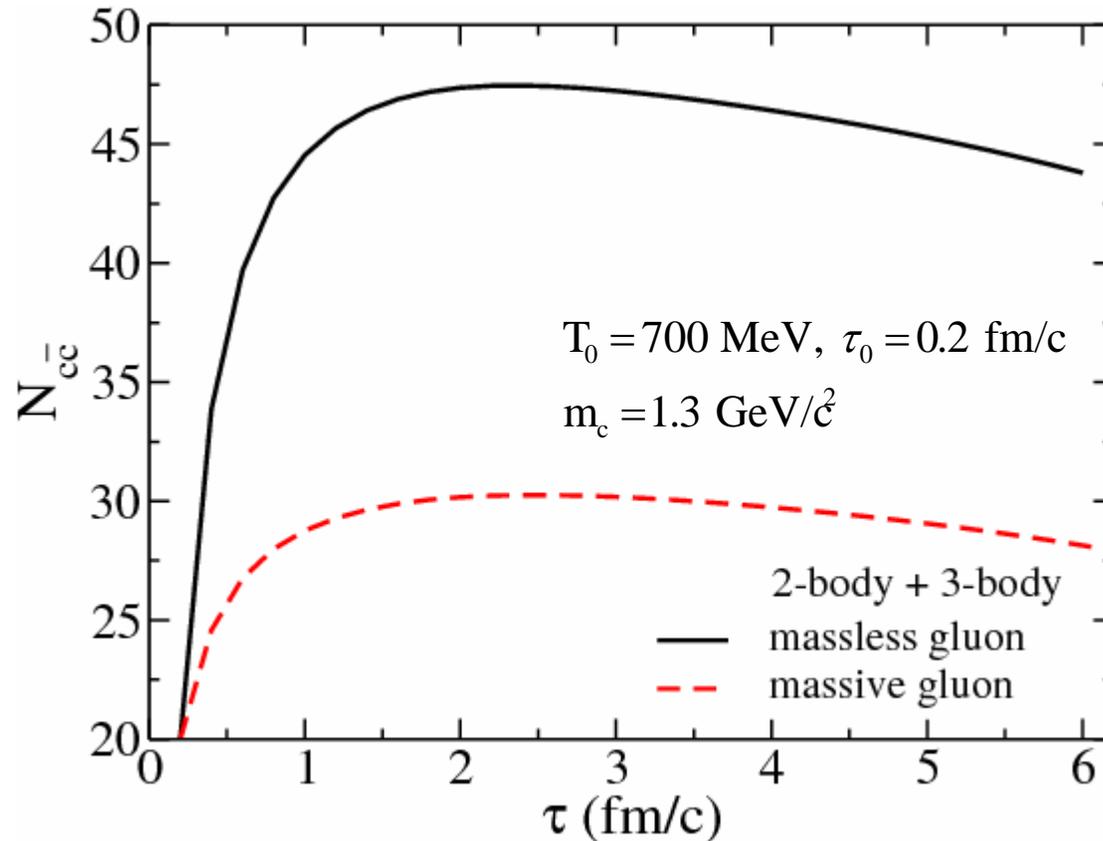
$$R \propto \frac{1}{3} \int \prod_{i=1}^5 d^3 p_i f_i(p_i) |M_{ggg \rightarrow c\bar{c}}|^2 \delta^{(4)}(p_1 + p_2 + p_3 - p_4 - p_5) \propto \langle \sigma_{c\bar{c} \rightarrow ggg} v \rangle n_c^{\text{eq}} n_{\bar{c}}^{\text{eq}}$$



Gluon density $\sim 0.5/\text{fm}^3$ at T_c
and much larger initially

- Negligible rate for massive gluons as the threshold becomes larger than the charm pair mass
- With massless gluons, the rate is comparable to that of two-body processes

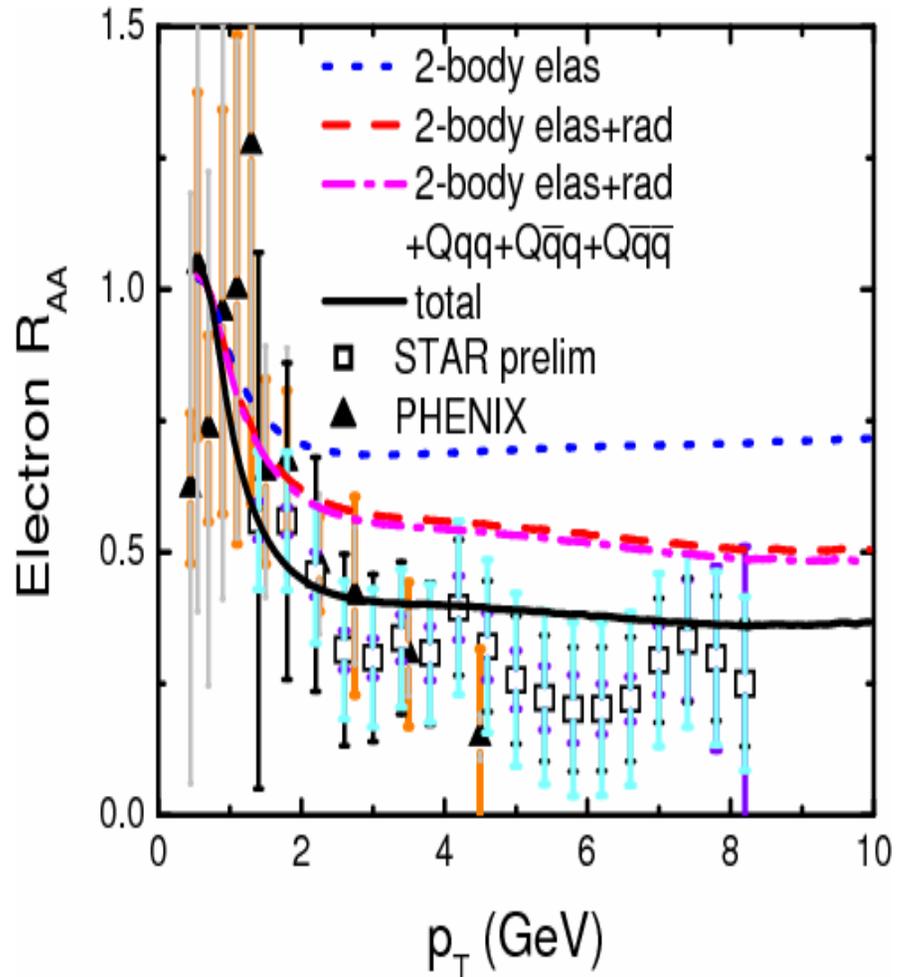
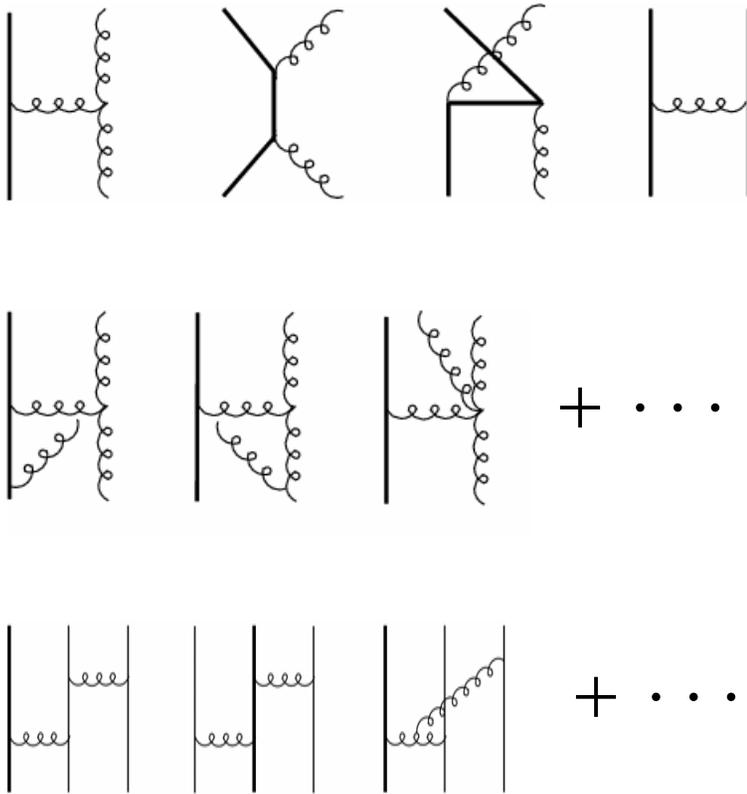
Time evolution of charm quark pairs at LHC including both two- and three-body interactions



Significant thermal production of charms from QGP of massless gluons

Nuclear modification factor for electrons from heavy meson decays

W. Liu & CMK, NPA 783, 233 (07); nucl-th/0603004



Reasonable agreement with data from Au+Au @ 200A GeV after including heavy quark three-body scattering.

Summary

- Thermal charm production rate increases \sim exponentially with the temperature of QGP.
- Next-to-leading order enhances thermal production rate by more than a factor of 2.
- Charm production from three-gluon interactions is important if gluons are massless.
- Thermal charm production could be important at LHC.
- Understanding thermal charm quark production is important for understanding charmonium production in HIC.
- LHC provides the possibility to search for charmed exotics such as charmed tetraquark mesons and pentaquark baryons.